

Recent developments in Bogoliubov Many-Body Perturbation Theory

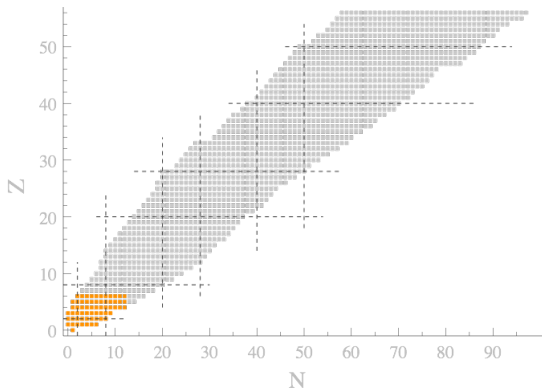
Pierre Arthuis

Unraveling the complexity of nuclear systems
ECT* - February 6th 2017

- ① A brief reminder on ab initio methods
- ② On Bogoliubov Many-Body Perturbation Theory

① A brief reminder on ab initio methods

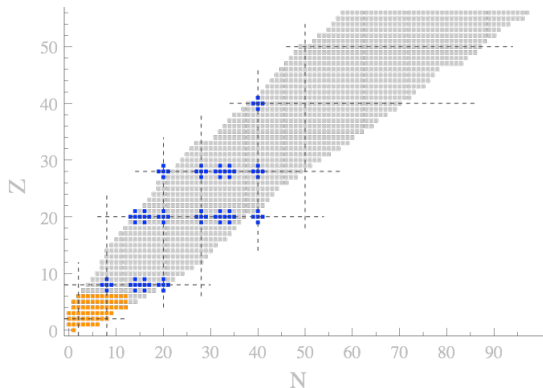
② On Bogoliubov Many-Body Perturbation Theory



Courtesy of V. Soma, T. Duguet

"Exact" *ab initio* methods

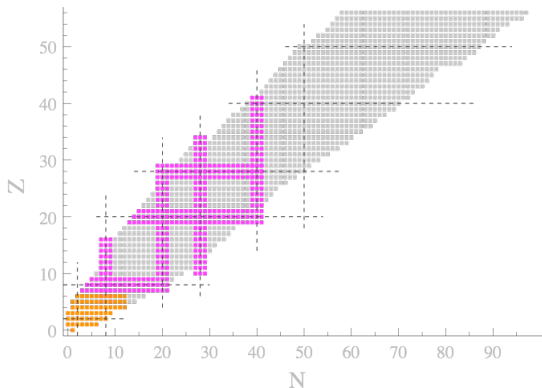
- Since the 80's
- GFMC, NCSM, FY



Courtesy of V. Soma, T. Duguet

Ab initio approaches for closed-shell nuclei

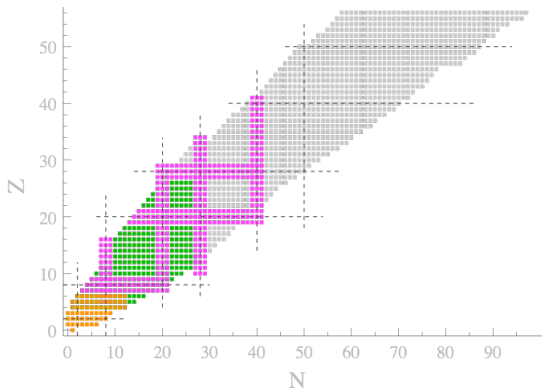
- Since the 2000's
- DSCGF, CC, IMSRG



Courtesy of V. Soma, T. Duguet

Non-perturbative *ab initio* approaches for open-shell nuclei

- Since the 2010's
- GSCGF, BCC, MR-IMSRG



Courtesy of V. Soma, T. Duguet

Ab initio shell model

- Since 2014
- Effective interaction via CC/IMSRG

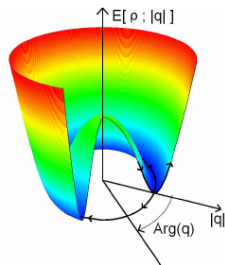
- 1 Consider point-like nucleons as appropriate degrees of freedom
- 2 Use interactions rooted in underlying theory (i.e. QCD)
- 3 Expand the many-body Schrödinger equation systematically
- 4 Truncate at a given order and solve using computational methods
- 5 Estimate systematic error

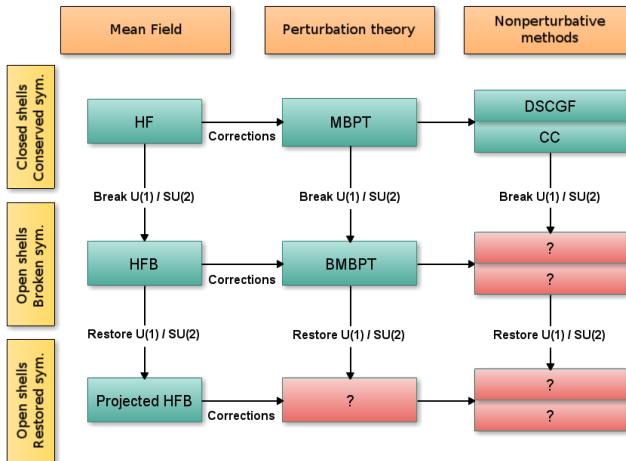
Symmetry breaking helps incorporating non-dynamical correlations:

- Superfluid character: $U(1)$ (particle number)
- Deformations: $SU(2)$ (angular momentum)

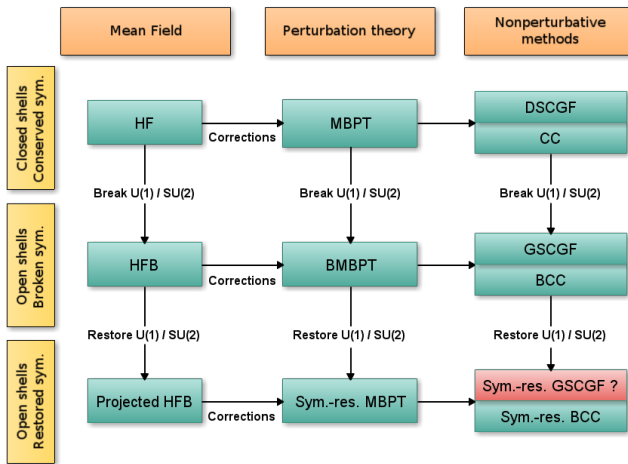
But nuclei carry good quantum numbers (e.g. number of particles)

⇒ Symmetries must eventually be restored





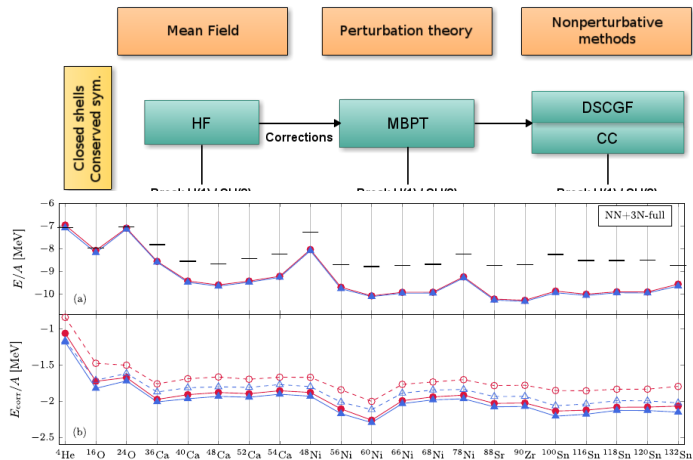
Expansion methods around unperturbed product state



MBPT: Recently (re)implemented with SRG-evolved H [Tichai *et al.* 2016]

GSCGF, BCC: Recently proposed and implemented [Somà *et al.* 2011, Signoracci *et al.* 2014]

Sym.-res. BCC & sym.-res. BMBPT: Recently proposed [Duguet 2015, Duguet & Signoracci 2016]



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- ① A brief reminder on ab initio methods
- ② On Bogoliubov Many-Body Perturbation Theory

- 1 Use a Bogoliubov vacuum $|\Phi\rangle$ with $\beta_k|\Phi\rangle = 0$ for all k
- 2 Define grand potential operator Ω from chiral interaction

$$\Omega \equiv H - \lambda A$$

then normal-order and split: $\Omega = \Omega_0 + \Omega_1$

- 3 Define evolved state in imaginary time

$$|\Psi(\tau)\rangle \equiv \mathcal{U}(\tau)|\Phi\rangle = e^{-\tau\Omega_0} \mathcal{T} e^{-\int_0^\tau d\tau \Omega_1(\tau)} |\Phi\rangle$$

- 4 Expand and truncate the grand potential kernel $\Omega(\tau) \equiv \langle \Psi(\tau) | \Omega | \Phi \rangle \dots$
...and the norm kernel $N(\tau) \equiv \langle \Psi(\tau) | \Phi \rangle$
- 5 Extract ground state energy via

$$E_0 = \lim_{\tau \rightarrow \infty} \frac{\Omega(\tau)}{N(\tau)} = \lim_{\tau \rightarrow \infty} \omega(\tau)$$

Inserting the operator Ω at time 0 and expanding

$$\begin{aligned} E_0 &= \lim_{\tau \rightarrow \infty} \frac{\langle \Psi(\tau) | \Omega | \Phi \rangle}{\langle \Psi(\tau) | \Phi \rangle} \\ &= \langle \Phi | \left\{ \Omega(0) - \int_0^\infty d\tau_1 \mathbb{T} [\Omega_1(\tau_1) \Omega(0)] \right. \\ &\quad \left. + \frac{1}{2!} \int_0^\infty d\tau_1 d\tau_2 \mathbb{T} [\Omega_1(\tau_1) \Omega_1(\tau_2) \Omega(0)] + \dots \right\} | \Phi \rangle_c \end{aligned}$$

Then expressing the grand potential in the qp basis

$$\Omega = \Omega^{00} + \frac{1}{1!} \sum_{k_1 k_2} \Omega_{k_1 k_2}^{11} \beta_{k_1}^\dagger \beta_{k_2} + \frac{1}{2!} \sum_{k_1 k_2} \left\{ \Omega_{k_1 k_2}^{20} \beta_{k_1}^\dagger \beta_{k_2}^\dagger + \Omega_{k_1 k_2}^{02} \beta_{k_2} \beta_{k_1} \right\} + \dots$$

$$\begin{aligned}
 E_0 = & \sum_{p=0}^{\infty} \frac{(-1)^p}{p!} \sum_{\substack{i_0+j_0=2,4 \\ \vdots \\ i_p+j_p=2,4}} \int_0^{\infty} d\tau_1 \dots d\tau_p \\
 & \times \sum_{\substack{k_1 \dots k_{i_1} \\ k_{i_1+1} \dots k_{i_1+j_1} \\ \vdots \\ l_1 \dots l_{i_p} \\ l_{i_p+1} \dots l_{i_p+j_p}}} \frac{\Omega_{k_1 \dots k_{i_1} k_{i_1+1} \dots k_{i_1+j_1}}^{i_1 j_1}}{(i_1)! (j_1)!} \dots \frac{\Omega_{l_1 \dots l_{i_p} l_{i_p+1} \dots l_{i_p+j_p}}^{i_p j_p}}{(i_p)! (j_p)!} \frac{\Omega_{m_1 \dots m_{i_0} m_{i_0+1} \dots m_{i_0+j_0}}^{i_0 j_0}}{(i_0)! (j_0)!} \\
 & \times \langle \Phi | T \left[\beta_{k_1}^{\dagger}(\tau_1) \dots \beta_{k_{i_1}}^{\dagger}(\tau_1) \beta_{k_{i_1+j_1}}(\tau_1) \dots \beta_{k_{i_1+1}}(\tau_1) \dots \right. \\
 & \quad \dots \beta_{l_1}^{\dagger}(\tau_p) \dots \beta_{l_{i_p}}^{\dagger}(\tau_p) \beta_{l_{i_p+j_p}}(\tau_p) \dots \beta_{l_{i_p+1}}(\tau_p) \\
 & \quad \left. \times \beta_{m_1}^{\dagger}(0) \dots \beta_{m_{i_0}}^{\dagger}(0) \beta_{m_{i_0+j_0}}(0) \dots \beta_{m_{i_0+1}}(0) \right] | \Phi \rangle_c
 \end{aligned}$$

All contributions computable algebraically and diagrammatically

Zero(first)- and first(second)-order diagrams

Diagrammatic representation of the grand potential Ω

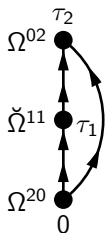
$$\Omega = \begin{array}{c} \bullet \\ \Omega^{00} \end{array} + \begin{array}{c} \uparrow \\ \bullet \\ \uparrow \\ \Omega^{11} \end{array} + \begin{array}{c} \swarrow \quad \nearrow \\ \bullet \\ \Omega^{20} \end{array} + \begin{array}{c} \nearrow \quad \swarrow \\ \bullet \\ \Omega^{02} \end{array} + \dots$$

Extracting and applying diagrammatic rules

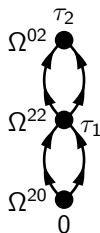
$$E_0^{(0+1)} = \begin{array}{c} \bullet \\ \Omega^{00} \end{array} \quad \begin{array}{c} \tau_1 \Omega^{02} \\ \bullet \\ \bullet \\ \Omega^{20} \end{array} \quad \begin{array}{c} \tau_1 \Omega^{04} \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ \Omega^{40} \end{array}$$

PE0.1 PE1.1 PE1.2

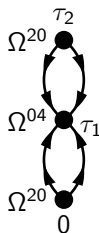
Second(third)-order diagrams



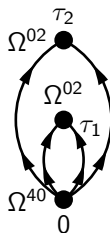
PE2.1



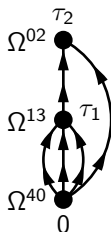
PE2.2



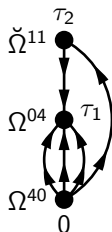
PE2.3



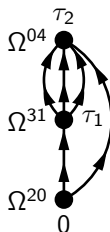
PE2.4



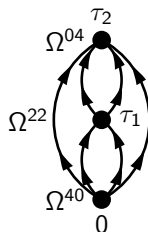
PE2.5



PE2.6



PE2.7



PE2.8

Derivation of all diagrams up to second(third) order

BMBPT must match standard MBPT in Slater determinant limit

→ Matching must be true at each order

→ Proof of consistent formalism for BMBPT

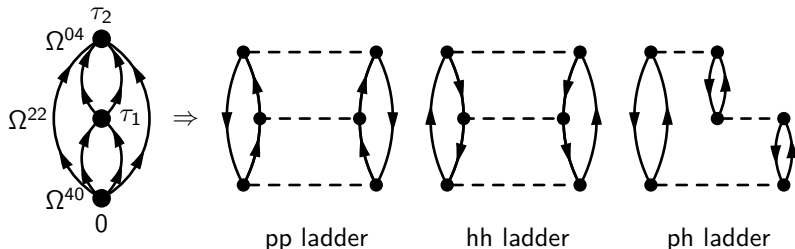
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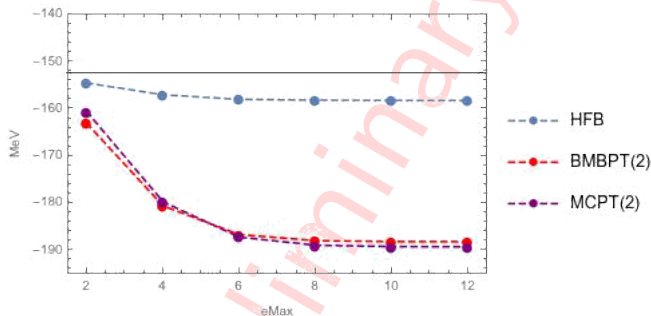
→ Proof of consistent formalism for BMBPT

BMBPT(3) diagrams match MBPT(3) ones exactly

Canonical HF-MBPT diagrams were recovered from PE2.8



First BMBPT(2) proof of principle calculation of ^{20}O :



[Arthuis, Tichai, Hergert, Roth, Duguet, in prep.]
[Tichai, Gebrerufael, Roth, in prep.]

using NN SRG-evolved chiral interaction (EM with $\Lambda = 500\text{MeV}$)

On MCPT:

- Multi-configurational MBPT
- Alternative method for open-shell nuclei

- Study of the complete O and Ca isotopic chains
 - First study using BMBPT
 - Comparison with other *ab initio* methods
- Push BMBPT to heavier nuclei
 - Can go further than other *ab initio* methods
 - Good test for the computational cost
- Implement particle-number restored BMBPT for the first time
 - Required for precise study of open-shell nuclei
 - Proof of concept of symmetry-restored BMBPT / BCC
- *Ab initio* driven EDF method [T. Duguet et al. 2015]
 - Safe/correlated/improvable off-diagonal EDF kernels
 - Based on PNR-BMBPT

- *Ab initio* methods are a powerful framework to study nuclei
 - ✓ Rigorous approach to the many-body problem
 - ✗ Computationally intensive
 - ✗ Cannot describe the whole nuclear chart
- Many-Body Perturbation Theory and its daughters are one of them
 - ✓ Computationally friendly
 - ✗ Potentially not as precise as others
- BMBPT has been formulated and is being implemented
 - ✓ First derivation and calculations up to second(third) order
 - ✓ Appropriate framework to tackle open-shell nuclei
 - ✓ Systematic studies at second(third) order to come
- Symmetry-restored BMBPT is the next step

BMBPT Project



P. Arthuis
T. Duguet
J.-P. Ebran



A. Tichai
R. Roth



H. Hergert

On broader aspects



M. Drissi
J. Ripoche



R. Lasserri