

# Coulomb phase stability and quark confinement

M. Asorey & A. Santagata

Universidad de Zaragoza



QCD-TNT-III

ECT\*, Trento, September 2013

# Analytic Flashes of QCD

- Perturbation theory: **Asymptotic freedom**
- Strong coupling: **Confinement**
- Intermediate regime: **Quarkonia (NRQCD)**

# Analytic Flashes of QCD

- Perturbation theory: **Asymptotic freedom**
- Strong coupling: **Confinement**
- Intermediate regime: **Quarkonia (NRQCD)**

## Pictures Confinement

- Lattice simulations
- Dual Superconductor
- AdS/CFT dualities

# Gribov's Quark Confinement

## One Electron Atomic Spectrum

$$E_{n,j} = mc^2 \left[ 1 + \left( \frac{Z\alpha}{n - (j + \frac{1}{2}) + \sqrt{(j + \frac{1}{2})^2 - (Z\alpha)^2}} \right)^2 \right]^{-\frac{1}{2}}$$

# Gribov's Quark Confinement

## One Electron Atomic Spectrum

$$E_{n,j} = mc^2 \left[ 1 + \left( \frac{Z\alpha}{n - (j + \frac{1}{2}) + \sqrt{(j + \frac{1}{2})^2 - (Z\alpha)^2}} \right)^2 \right]^{-\frac{1}{2}}$$

Vacuum instability for

$$Z > 137$$

The running coupling of  $\alpha_s$  in the IR limit is large !!.

# Gribov's Quark Confinement

Quark Green function Dyson-Schwinger equation

$$D^2 G^{-1}(k) = \frac{(N^2 - 1)\alpha_s}{2N\pi} \partial_\mu G^{-1}(k) G(k) \partial^\mu G^{-1}(k) + \mathcal{O}\left(\frac{\alpha^2}{2}\right)$$

# Gribov's Quark Confinement

Quark Green function Dyson-Schwinger equation

$$D^2 G^{-1}(k) = \frac{(N^2 - 1)\alpha_s}{2N\pi} \partial_\mu G^{-1}(k) G(k) \partial^\mu G^{-1}(k) + \mathcal{O}\left(\frac{\alpha^2}{2}\right)$$

Instability for large strong coupling  $\alpha_s$  constant

$$\alpha_s > \alpha_{crit} = \frac{2N\pi}{(N^2 - 1)} \left( 1 - \sqrt{\frac{2}{3}} \right)$$

$$\alpha_{crit} = 0.43 \quad N = 3$$

# Gribov's Quark Confinement

Heavy quark potential:

$$S_E^{YM}(A) = -\frac{1}{2g_s^2} \int d^4x \operatorname{Tr}(F^{\mu\nu} F_{\mu\nu}) + Q \int dx^0 A_0^3$$



# Gribov's Quark Confinement

Heavy quark potential:

$$S_E^{YM}(A) = -\frac{1}{2g_s^2} \int d^4x \operatorname{Tr}(F^{\mu\nu}F_{\mu\nu}) + Q \int dx^0 A_0^3$$

Solution of motion equations (Coulomb)

$$(A_0)^3(\vec{x}) = i \frac{g_s^2 Q}{4\pi |\vec{x}|} = i \frac{\alpha}{|\vec{x}|}, \quad \alpha = \frac{g_s^2 Q}{4\pi},$$

# Gribov's Quark Confinement

Heavy quark potential:

$$S_E^{YM}(A) = -\frac{1}{2g_s^2} \int d^4x \operatorname{Tr}(F^{\mu\nu}F_{\mu\nu}) + Q \int dx^0 A_0^3$$

Solution of motion equations (Coulomb)

$$(A_0)^3(\vec{x}) = i \frac{g_s^2 Q}{4\pi |\vec{x}|} = i \frac{\alpha}{|\vec{x}|}, \quad \alpha = \frac{g_s^2 Q}{4\pi},$$

Instability of Euclidean functional integral

$$\delta^{(2)} S = - \int d^4x \operatorname{Tr} \left( \tau^\mu (-\delta_{\mu\nu} D^2 + D_\mu D_\nu - 2[F_{\mu\nu} \cdot]) \tau^\nu \right),$$

# Gribov's Unstability

Eigenvalue problem

$$(-\delta_{\mu\nu}D^2 + D_\mu D_\nu - 2[F_{\mu\nu}, \cdot])\tau^\nu = \lambda^2 \tau_\mu$$

# Gribov's Unstability

Eigenvalue problem

$$(-\delta_{\mu\nu}D^2 + D_\mu D_\nu - 2[F_{\mu\nu}, \cdot])\tau^\nu = \lambda^2 \tau_\mu$$

Unstable magnetic modes

$$\vec{\tau}(\vec{x}) = \frac{\vec{x} \times \vec{n}}{|\vec{x}|} \phi(|\vec{x}|) \mathbf{T}_{12}, \quad \tau_0 = 0 \quad (l = 1),$$

$$\left( \frac{d^2}{dr^2} + \frac{2}{r} \frac{d}{dr} - \frac{2 - \alpha^2}{r^2} \right) \phi(r) = \lambda^2 \phi(r)$$

# Singular Hamiltonians

One dimensional problem

$$H = -\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} \frac{g^2}{x^2}$$

in the half line  $x \in (0, \infty)$

# Singular Hamiltonians

One dimensional problem

$$H = -\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} \frac{g^2}{x^2}$$

in the half line  $x \in (0, \infty)$

## Applications

- 2-body problems. Centrifugal potential
- 3-body problem (**Efimov effect**)
- Relativistic Coulomb problem

# Singular Hamiltonians

One dimensional problem

$$H = -\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} \frac{g^2}{x^2}$$

in the half line  $x \in (0, \infty)$

## Applications

- 2-body problems. Centrifugal potential
- 3-body problem (**Efimov effect**)
- Relativistic Coulomb problem
- **Gribov's approach to confinement**

# Singular Hamiltonians

One dimensional problem

$$H = -\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} \frac{g^2}{x^2}$$

in the half line  $x \in (0, \infty)$

## Applications

- 2-body problems. Centrifugal potential
- 3-body problem (**Efimov effect**)
- Relativistic Coulomb problem
- $H$  is symmetric but in this case we have **three different regimes**



# Singular Hamiltonians

i)  $\frac{3}{4} < g^2 \Rightarrow$  unique selfadjoint extension

# Singular Hamiltonians

i)  $\frac{3}{4} < g^2 \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < g^2 < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\Lambda \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} \left( 1 + 2\sqrt{\frac{1}{4} + g^2} \tanh\left[\sqrt{\frac{1}{4} + g^2} \log(\Lambda x)\right] \right) \psi(x),$$

# Singular Hamiltonians

i)  $\frac{3}{4} < g^2 \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < g^2 < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\Lambda \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} \left( 1 + 2\sqrt{\frac{1}{4} + g^2} \tanh[\sqrt{\frac{1}{4} + g^2} \log(\Lambda x)] \right) \psi(x),$$

iii)  $g^2 < -\frac{1}{4} \Rightarrow$  family of bc parametrized by  $\Lambda \in \mathbb{R}$

$$\lim_{x=0} \left( 1 + 2\sqrt{-g^2 - \frac{1}{4}} \cot[\sqrt{-g^2 - \frac{1}{4}} \log(\Lambda x)] \right) \psi(x) = \lim_{x=0} 2x\psi'(x)$$

with periodicity in  $\Lambda$

$$\Lambda \Leftrightarrow \Lambda e^{2\pi/\sqrt{-g^2-1/4}}$$

Efimov effect

# Singular Hamiltonians

Infinite sequence of negative energy levels

$$E_n = 2\Lambda^2 \exp \left( \frac{-2\pi n}{2\sqrt{-g^2-1/4}} + \frac{i}{2\sqrt{-g^2-1/4}} \log \frac{\Gamma(1-i\sqrt{-g^2-1/4})}{\Gamma(1+i\sqrt{-g^2-1/4})} \right)$$

$$\frac{E_{n+1}}{E_n} = \exp[-2\pi/(\sqrt{g-1/4})].$$

Partial breaking of conformal symmetry (**Conformal Anomaly**)

$$G = \frac{1}{4}(xp + px)$$

$$D = tH - G$$

$$K = -t^2 + 2Dt + \frac{1}{2}x^2$$

# Quark Coulomb Instability

- Zero modes with  $j = 0$  (pure gauge modes)

# Quark Coulomb Instability

- Zero modes with  $j = 0$  (pure gauge modes)
- Three different unstable regimes  $j = 1$ 
  - i)  $\alpha^2 < \frac{5}{4} \Rightarrow$  no negative eigenvalues  
(Stability of Coulomb potential)
  - ii)  $\frac{5}{4} < \alpha^2 < \frac{9}{4} \Rightarrow$  one negative eigenvalue  
(Instability of Coulomb potential)
  - iii)  $\alpha^2 > \frac{9}{4} \Rightarrow \infty$ -negative eigenvalues  
(Instability of Coulomb potential)

$\alpha^2$

$\frac{5}{4}$

$\frac{9}{4}$



Broken Conformal Symmetry  $\Lambda$

M.A. & A. Santagata (2010)

# Quark-Antiquark: Meson

Coulomb potential

$$(A_0)^3(\vec{x}) = \frac{i\alpha}{|\vec{x} - L\vec{e}_3|} - \frac{i\alpha}{|\vec{x} + L\vec{e}_3|},$$

# Quark-Antiquark: Meson

Coulomb potential

$$(A_0)^3(\vec{x}) = \frac{i\alpha}{|\vec{x} - L\vec{e}_3|} - \frac{i\alpha}{|\vec{x} + L\vec{e}_3|},$$

Unstable magnetic modes

$$\vec{\tau}(\vec{x}) = \frac{\vec{x} \times \vec{e}_3}{\rho} \phi(\rho, z) T_{12}, \quad \tau_0 = 0 \quad (m = 1),$$



# Quark-Antiquark: Meson

Coulomb potential

$$(A_0)^3(\vec{x}) = \frac{i\alpha}{|\vec{x} - L\vec{e}_3|} - \frac{i\alpha}{|\vec{x} + L\vec{e}_3|},$$

Unstable magnetic modes

$$\vec{\tau}(\vec{x}) = \frac{\vec{x} \times \vec{e}_3}{\rho} \phi(\rho, z) \mathbb{T}_{12}, \quad \tau_0 = 0 \quad (m = 1),$$

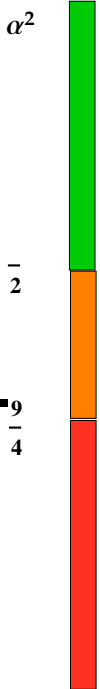
$$\left[ \frac{\partial^2}{\partial \rho^2} + \frac{\partial^2}{\partial z^2} - \frac{3}{4\rho^2} + \left( \frac{\alpha}{\sqrt{\rho^2 + (z - L)^2}} - \frac{\alpha}{\sqrt{\rho^2 + (z + L)^2}} \right)^2 \right] \phi(\rho, z) = \lambda^2 \phi(\rho, z)$$

# Quark-Antiquark Coulomb Instability

- Zero modes with  $j = 0$  (pure gauge modes)

# Quark-Antiquark Coulomb Instability

- Zero modes with  $j = 0$  (pure gauge modes)
- Three different unstable regimes  $j = 1$ 
  - i)  $\alpha^2 < 2 \Rightarrow$  no negative eigenvalues  
(Stability of Coulomb potential)
  - ii)  $2 < \alpha^2 < \frac{9}{4} \Rightarrow$  two negative eigenvalues at large distances and none at short distances.  
(Instability of Coulomb potential)
  - iii)  $\alpha^2 > \frac{9}{4} \Rightarrow \infty$ -negative eigenvalues  
(Instability of Coulomb potential)

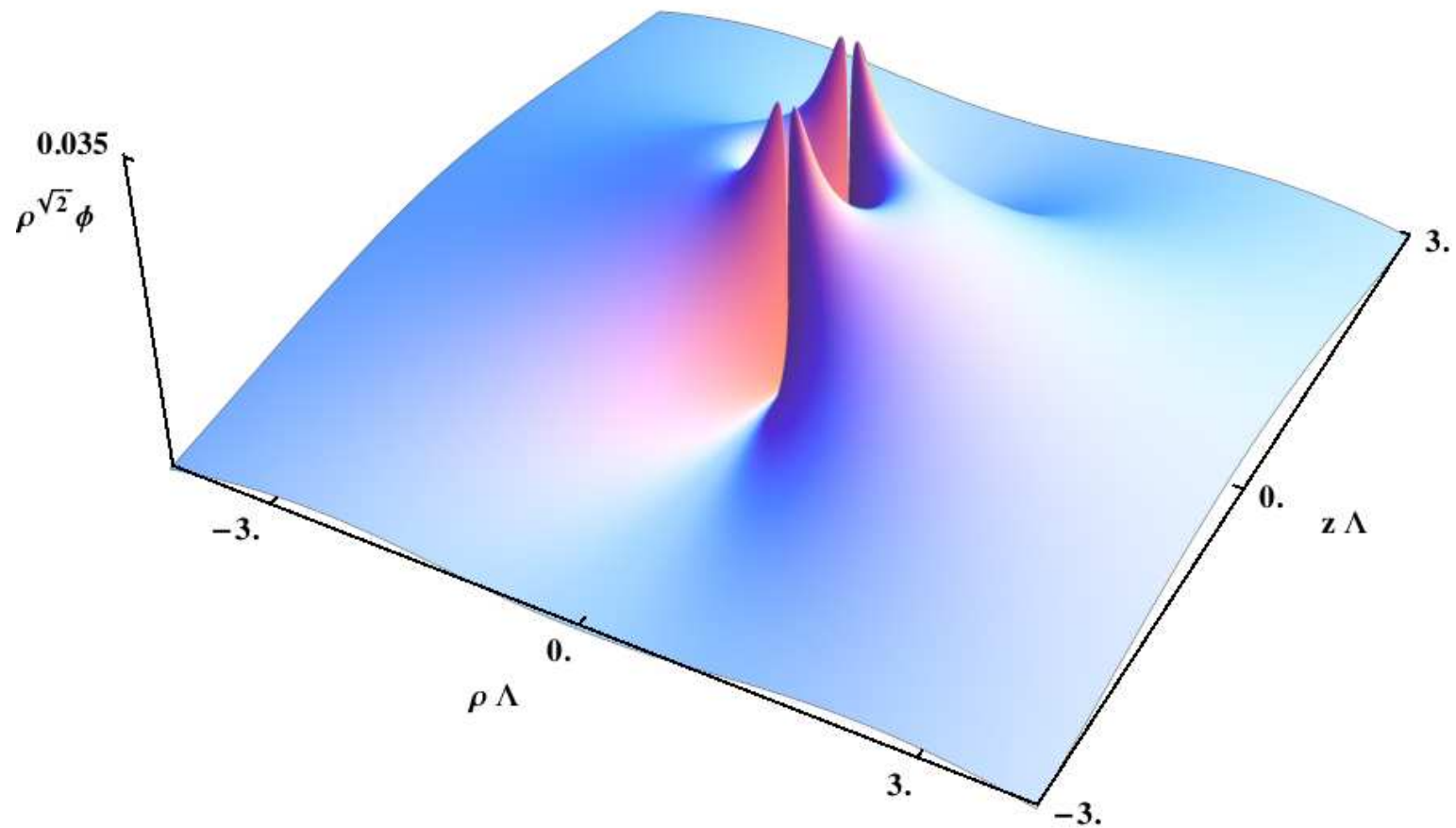


Broken Conformal Symmetry  $\Lambda$   
M.A. & A. Santagata (2012)

# Unstable solutions

$$2 < \alpha^2 < \frac{9}{4}$$

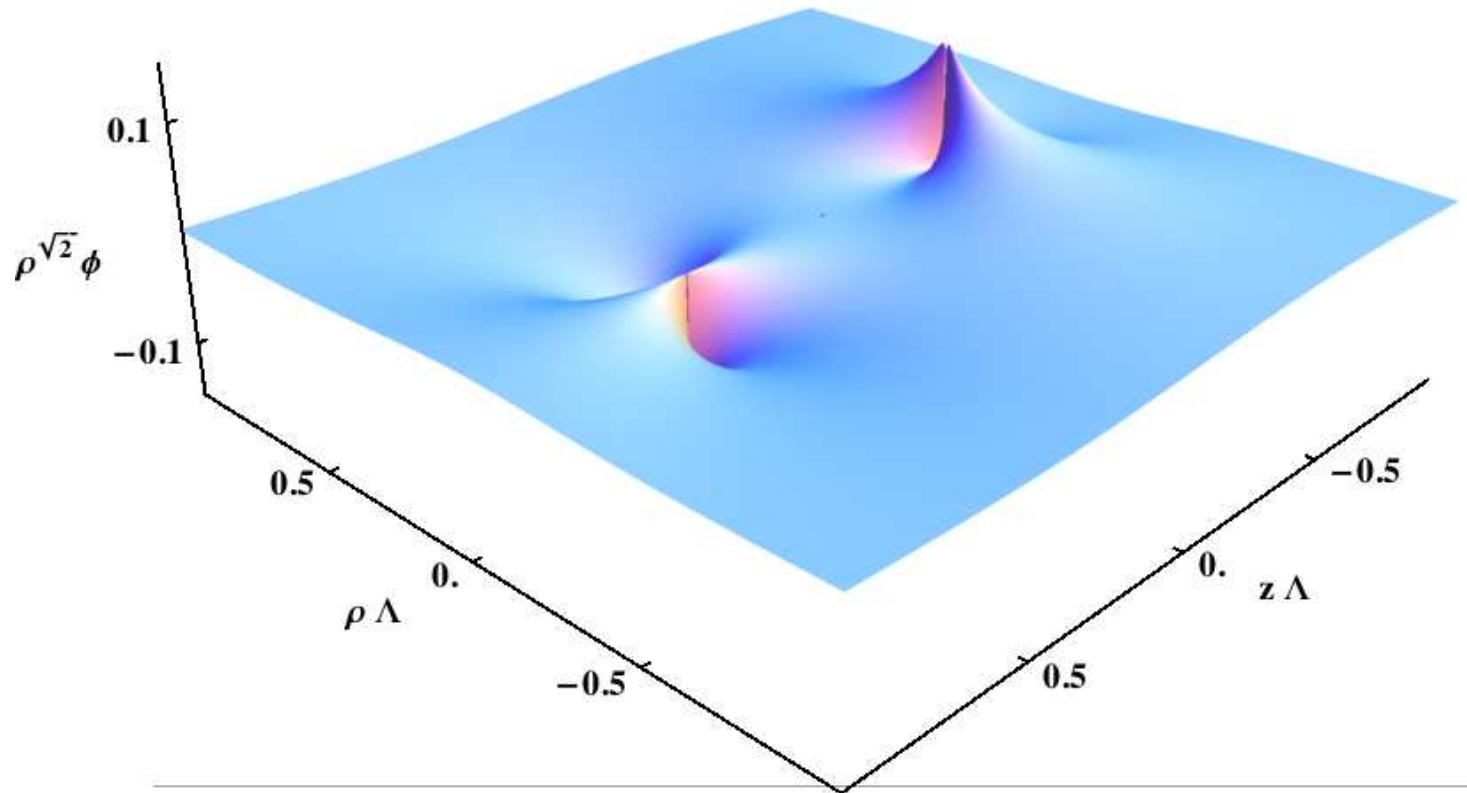
Symmetric solution (Thick string)



# Unstable solutions

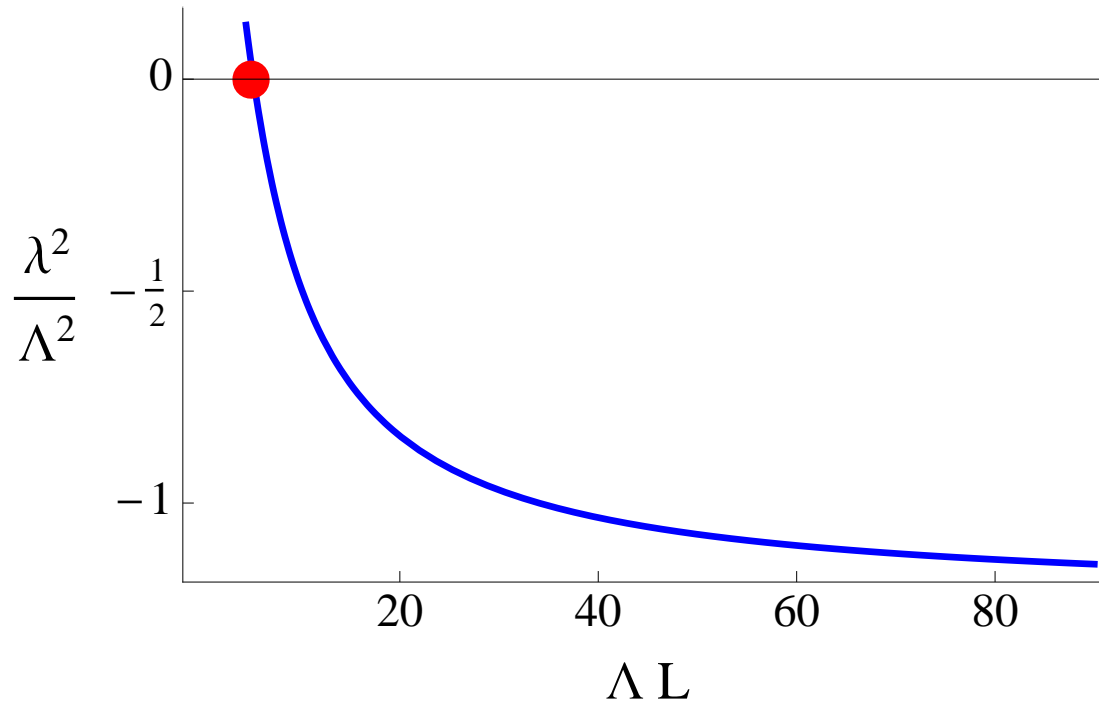
$$2 < \alpha^2 < \frac{9}{4}$$

Anti-symmetric solution



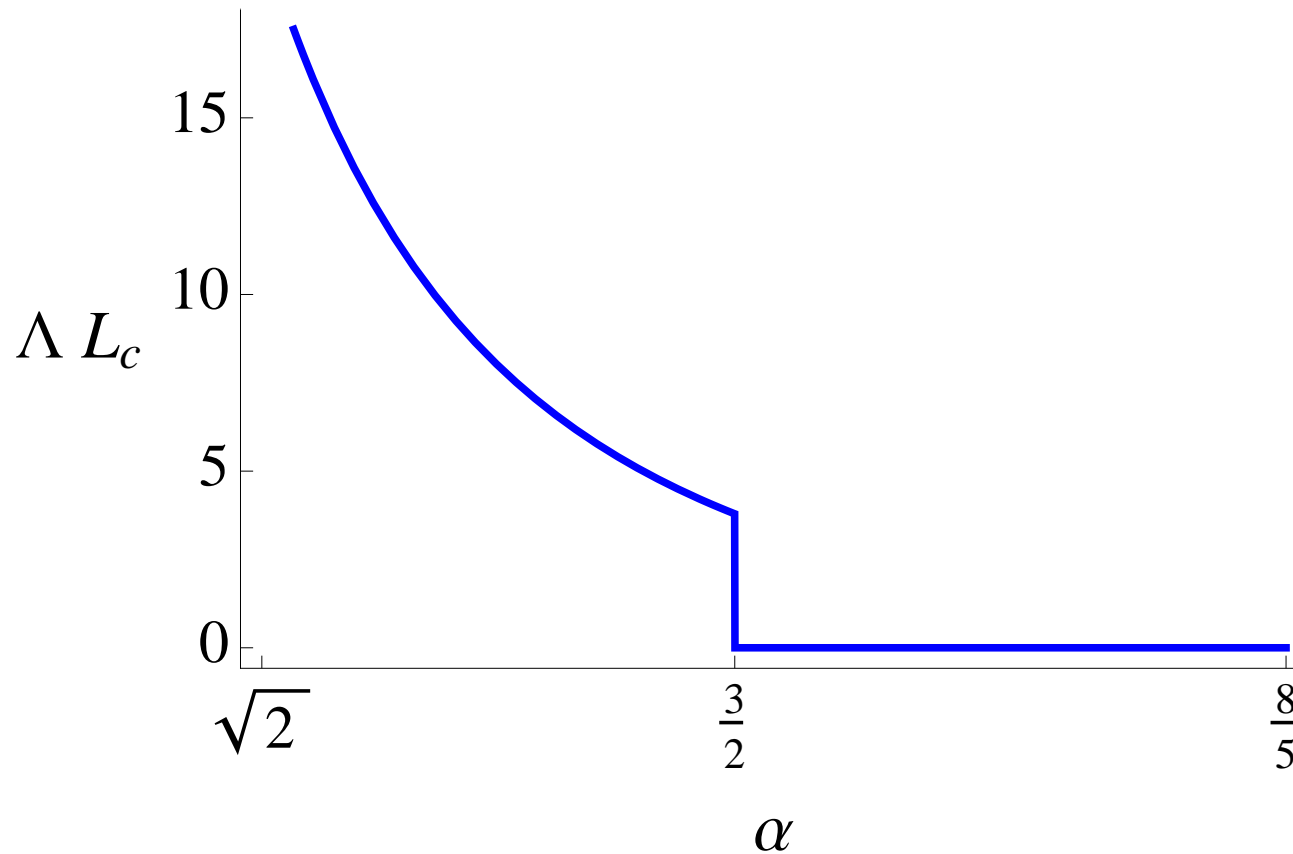
# Existence of a critical $q - \bar{q}$ distance

Unstable modes disappear at short  $q - \bar{q}$  distances



# Existence of a critical $q - \bar{q}$ distance

Dependence of critical  $q - \bar{q}$  distance on  $\alpha$



$$\alpha = \sqrt{2}$$

$$\alpha_G = 0.43$$

# COULOMB PHASE INSTABILITIES

- Gribov picture of confinement derived from first principles



# COULOMB PHASE INSTABILITIES

- Gribov picture of confinement derived from first principles
- **Weak Coupling regime**  $\alpha^2 < 2$ :  
Coulomb phase is stable (perturbative regime)

# COULOMB PHASE INSTABILITIES

- Gribov picture of confinement derived from first principles
- **Weak Coupling regime**  $\alpha^2 < 2$ :  
Coulomb phase is stable (perturbative regime)
- **Strong Coupling regime**  $\alpha^2 > \frac{9}{4}$ :  
Coulomb phase is unstable (confinement)

# COULOMB PHASE INSTABILITIES

- Gribov picture of confinement derived from first principles
- **Weak Coupling regime**  $\alpha^2 < 2$ :  
Coulomb phase is stable (perturbative regime)
- **Strong Coupling regime**  $\alpha^2 > \frac{9}{4}$ :  
Coulomb phase is unstable (confinement)
- **Intermediate regime**  $2 < \alpha^2 < \frac{9}{4}$ :  
there is a critical quark distance  $L_c$

# COULOMB PHASE INSTABILITIES

- Gribov picture of confinement derived from first principles
- **Weak Coupling regime**  $\alpha^2 < 2$ :  
Coulomb phase is stable (perturbative regime)
- **Strong Coupling regime**  $\alpha^2 > \frac{9}{4}$ :  
Coulomb phase is unstable (confinement)
- **Intermediate regime**  $2 < \alpha^2 < \frac{9}{4}$ :  
there is a critical quark distance  $L_c$ 
  - $L < L_c$  Coulomb phase stable (**asymt. freedom**)
  - $L > L_c$  Coulomb phase unstable (**confinement**)

# COULOMB PHASE INSTABILITIES

- Stationary phase approximation points out Coulomb phase instabilities

# COULOMB PHASE INSTABILITIES

- Stationary phase approximation points out Coulomb phase instabilities
- Reflection positivity is broken in Coulomb background

# COULOMB PHASE INSTABILITIES

- Stationary phase approximation points out Coulomb phase instabilities
- Reflection positivity is broken in Coulomb background
- Compatibility static quark potentials and lattice results

# COULOMB PHASE INSTABILITIES

- Stationary phase approximation points out Coulomb phase instabilities
- Reflection positivity is broken in Coulomb background
- Compatibility static quark potentials and lattice results
- Gribov picture does not applies to SUSY theories.

## OUTLOOK



# COULOMB PHASE INSTABILITIES

- Stationary phase approximation points out Coulomb phase instabilities
- Reflection positivity is broken in Coulomb background
- Compatibility static quark potentials and lattice results
- Gribov picture does not applies to SUSY theories.

## OUTLOOK

- Finite temperature phase transition?
- Chiral symmetry breaking?

# Singular Hamiltonians and AdS/CFT

- Free scalar field in **anti-deSitter** space-time

$$S(\phi) = \frac{1}{2} \int \frac{dt dz}{z^2} \int d^3x \left[ |\dot{\phi}|^2 + |\partial_z \phi|^2 - |\nabla \phi|^2 - m^2 |\phi|^2 \right]$$

# Singular Hamiltonians and AdS/CFT

- Free scalar field in **anti-deSitter** space-time

$$S(\phi) = \frac{1}{2} \int \frac{dt dz}{z^2} \int d^3x \left[ |\dot{\phi}|^2 + |\partial_z \phi|^2 - |\nabla \phi|^2 - m^2 |\phi|^2 \right]$$

- Effective Hamiltonian

$$H = -\frac{1}{2} \frac{d^2}{dz^2} + \frac{m^2 + \frac{15}{4}}{z^2}$$

in the half line  $z \in (0, \infty)$

- $H$  is symmetric but in this case we have **three** different regimes

# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < m^2 + \frac{15}{4} < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} (1 + 2g \tanh[g \log(\alpha x)]) \psi(x),$$

with  $g = \sqrt{4 + m^2}$

# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < m^2 + \frac{15}{4} < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} (1 + 2g \tanh[g \log(\alpha x)]) \psi(x),$$

with  $g = \sqrt{4 + m^2}$

iii)  $m^2 + \frac{15}{4} < -\frac{1}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} (1 + 2g \cot[g \log(\alpha x)]) \psi(x) = \lim_{x=0} 2x\psi'(x)$$

with periodicity in  $\alpha$

$$\alpha \equiv \alpha e^{2\pi/g}$$

**Efimov effect**

# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < m^2 + \frac{15}{4} < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} (1 + 2g \tanh[g \log(\alpha x)]) \psi(x),$$

with  $g = \sqrt{4 + m^2}$



# Singular Hamiltonians and AdS/CFT

i)  $\frac{3}{4} < m^2 + \frac{15}{4} \Rightarrow$  unique selfadjoint extension

ii)  $-\frac{1}{4} < m^2 + \frac{15}{4} < \frac{3}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} 2x\psi'(x) = \lim_{x=0} (1 + 2g \tanh[g \log(\alpha x)]) \psi(x),$$

with  $g = \sqrt{4 + m^2}$

iii)  $m^2 + \frac{15}{4} < -\frac{1}{4} \Rightarrow$  family of bc parametrized by  $\alpha \in \mathbb{R}$

$$\lim_{x=0} (1 + 2g \cot[g \log(\alpha x)]) \psi(x) = \lim_{x=0} 2x\psi'(x)$$

with periodicity in  $\alpha$

$$\alpha \equiv \alpha e^{2\pi/g}$$

**Efimov effect**

# BACKUP SLIDES

$$x = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \cos \varphi$$

$$y = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \sin \varphi$$

$$z = a\mu\eta$$

$$\mu \geq 1 \quad -1 \leq \eta \leq 1$$

# BACKUP SLIDES

$$x = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \cos \varphi$$

$$y = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \sin \varphi$$

$$z = a\mu\eta$$

$$\mu \geq 1 \quad -1 \leq \eta \leq 1$$

# BACKUP SLIDES

$$x = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \cos \varphi$$

$$y = a\sqrt{(\mu^2 - 1)(1 - \eta^2)} \sin \varphi$$

$$z = a\mu\eta$$

$$\mu \geq 1 \quad -1 \leq \eta \leq 1$$