

ECT*

EUROPEAN CENTRE FOR THEORETICAL STUDIES
IN NUCLEAR PHYSICS AND RELATED AREAS

2nd Workshop on The Proton mass:
At the Heart of Most Visible Matter

3-7 April 2017

Longitudinal and transverse spin-orbit correlations

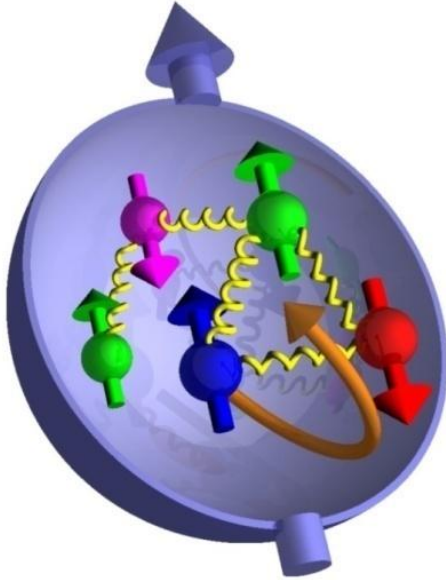
Based on : **[Bhooah, C.L. (2017)]**
[C.L. (2014)]

Cédric Lorcé



April 4, ECT*, Trento, Italy

Outline



- Energy-momentum tensor
- Spin-dependent decomposition
- Link with GPDs
- Lattice and quark model estimates
- Conclusions

Energy-momentum tensor

A lot of interesting physics is contained in the EM tensor !

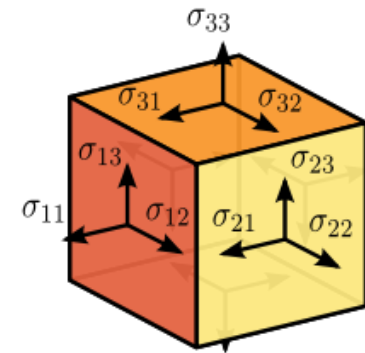
$$T^{\mu\nu} = \begin{bmatrix} \text{Energy density} & \text{Momentum density} & & \\ T^{00} & T^{01} & T^{02} & T^{03} \\ \text{Energy flux} & \text{Momentum flux} & & \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \end{bmatrix}$$

Shear stress
Normal stress (pressure)

Rest frame

$$M = \int d^3r T^{00}(\vec{r})$$

$$L^i = \int d^3r \epsilon^{ijk} r^j T^{0k}(\vec{r})$$



Energy-momentum tensor

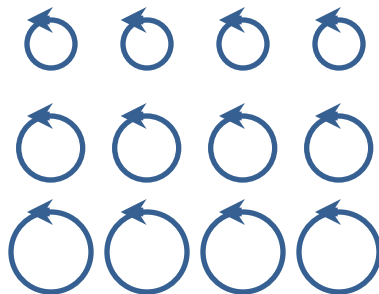
In presence of spin density

$$T^{0i} \neq T^{i0}$$

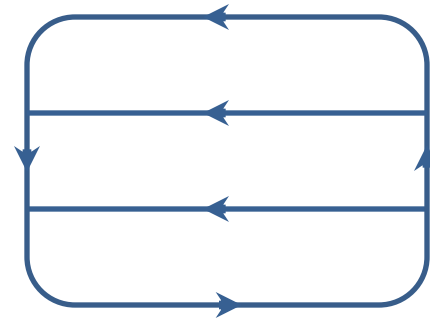
$$\left[\begin{array}{l} J^{\mu\alpha\beta} = x^\alpha T^{\mu\beta} - x^\beta T^{\mu\alpha} + S^{\mu\alpha\beta} \\ \partial_\mu T^{\mu\nu} = \partial_\mu J^{\mu\alpha\beta} = 0 \Rightarrow T^{[\alpha\beta]} = -\partial_\mu S^{\mu\alpha\beta} \end{array} \right.$$

Belinfante
« improvement »

$$\begin{aligned} T_B^{\mu\nu} &\equiv T^{\mu\nu} + \frac{1}{2} \partial_\lambda [S^{\lambda\mu\nu} + S^{\mu\nu\lambda} + S^{\nu\mu\lambda}] \\ &= T^{\nu\mu} \end{aligned}$$



Spin density gradient



Four-momentum circulation

Rest frame

$$M = \int d^3r T_B^{00}(\vec{r})$$

$$J^i = \int d^3r \epsilon^{ijk} r^j T_B^{0k}(\vec{r})$$

No « spin » contribution !

Energy-momentum tensor

Quark energy-momentum tensor

$$\hat{T}_q^{\mu\nu} = \bar{\psi} \gamma^\mu i \overleftrightarrow{D}^\nu \psi$$

$$\bar{\psi} \gamma^{[\mu} i \overleftrightarrow{D}^{\nu]} \psi = -\varepsilon^{\mu\nu\alpha\beta} \partial_\alpha (\bar{\psi} \gamma_\beta \gamma_5 \psi)$$

based on QCD EOM

General parametrization

[Bakker, Leader, Trueman (2004)]

$$\langle p' | \hat{T}^{\mu\nu} | p \rangle = \bar{u}(p') \left[\frac{P^{\{\mu} \gamma^{\nu\}}}{2} A(t) + \frac{P^{\{\mu} i \sigma^{\nu\} \Delta}}{4M} B(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} C(t) \right. \\ \left. + \underbrace{M g^{\mu\nu} \bar{C}(t)}_{\text{Non-conservation}} + \underbrace{\frac{P^{[\mu} \gamma^{\nu]}}{2} D(t)}_{\text{Asymmetry}} \right] u(p)$$

Higher twist

$$\begin{aligned} A_q + A_G &= 1 \\ B_q + B_G &= 0 \\ \bar{C}_q + \bar{C}_G &= 0 \end{aligned}$$

Sum rules

$$J_z = \frac{1}{2} [A(0) + B(0)]$$

[Ji (1997)]

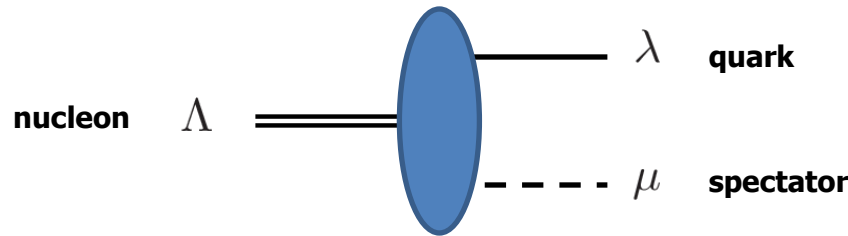
$$L_z = \frac{1}{2} [A(0) + B(0) + \underbrace{D(0)}_{-2S_z}]$$

[Shore, White (2000)]

t dependence \longleftrightarrow spatial distribution

Cf. Luca Mantovani's talk

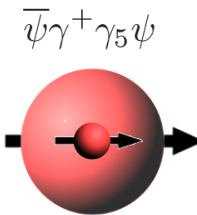
Proton spin structure



$$l_z = \Lambda - \lambda - \mu$$

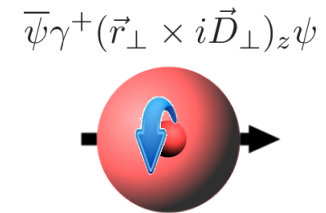
Quark spin

$$\langle\langle S_z^q \rangle\rangle \sim \frac{1}{2(2\pi)^3} \sum_{\Lambda, \lambda, \mu} \Lambda \lambda |\psi_{\lambda, \mu}^\Lambda|^2 \sim \langle S_z^N S_z^q \rangle$$



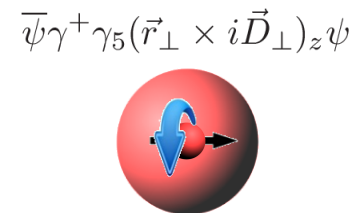
Quark OAM

$$\langle\langle L_z^q \rangle\rangle \sim \frac{1}{2(2\pi)^3} \sum_{\Lambda, \lambda, \mu} \Lambda l_z |\psi_{\lambda, \mu}^\Lambda|^2 \sim \langle S_z^N L_z^q \rangle$$



Quark spin-orbit correlation

$$\langle\langle C_z^q \rangle\rangle \sim \frac{1}{2(2\pi)^3} \sum_{\Lambda, \lambda, \mu} \lambda l_z |\psi_{\lambda, \mu}^\Lambda|^2 \sim \langle S_z^q L_z^q \rangle$$



Parity-odd energy-momentum tensor

Chiral decomposition

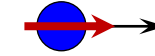
$$\hat{T}^{\mu\nu} = \bar{\psi} \gamma^\mu \frac{i}{2} \overleftrightarrow{D}^\nu \psi$$

$$\hat{T}_5^{\mu\nu} = \bar{\psi} \gamma^\mu \gamma_5 \frac{i}{2} \overleftrightarrow{D}^\nu \psi$$

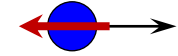
$$\hat{T}^{+\nu} = \hat{T}_R^{+\nu} + \hat{T}_L^{+\nu}$$

$$\hat{T}_5^{+\nu} = \hat{T}_R^{+\nu} - \hat{T}_L^{+\nu}$$

Right-handed



Left-handed



$$\hat{T}_a^{+\nu} = \bar{\psi}_a \gamma^+ \frac{i}{2} \overleftrightarrow{D}^\nu \psi_a \quad a = R, L$$

$$\psi_{R,L} = \frac{1 \pm \gamma_5}{2} \psi$$

General parametrization

[C.L. (2014)]

$$\langle p' | \hat{T}_5^{\mu\nu} | p \rangle = \bar{u}(p') \left[\frac{P^{\{\mu} \gamma^{\nu\}} \gamma_5}{2} \tilde{A}(t) + \frac{P^{\{\mu} \Delta^{\nu\}} \gamma_5}{4M} \tilde{B}(t) \right. \\ \left. + \frac{P^{[\mu} \gamma^{\nu]} \gamma_5}{2} \tilde{C}(t) + \frac{P^{[\mu} \Delta^{\nu]} \gamma_5}{4M} \tilde{D}(t) + Mi \sigma^{\mu\nu} \gamma_5 \tilde{F}(t) \right] u(p)$$

Higher twist

Longitudinal spin-orbit correlations

$$C_z = \frac{1}{2} [\tilde{A}(0) + \tilde{C}(0)] \quad \text{Asymmetric}$$


$$\mathbb{C}_z = \frac{1}{2} \tilde{A}(0) \quad \text{Belinfante}$$

Chiral-odd energy-momentum tensor

Transversity decomposition

$$\begin{aligned}\hat{T}^{\mu\nu} &= \bar{\psi} \gamma^\mu \frac{i}{2} \overleftrightarrow{D}^\nu \psi & \hat{T}^{+\nu} &= \hat{T}_\uparrow^{+\nu} + \hat{T}_\downarrow^{+\nu} \\ \hat{T}_5^{\lambda\mu\nu} &= \bar{\psi} i \sigma^{\lambda\mu} \gamma_5 \frac{i}{2} \overleftrightarrow{D}^\nu \psi & \hat{T}_5^{j+\nu} &= \hat{T}_\uparrow^{j+\nu} - \hat{T}_\downarrow^{j+\nu}\end{aligned}$$

Transversity basis



$$\begin{aligned}\hat{T}_a^{+\nu} &= \bar{\psi}_a \gamma^+ \frac{i}{2} \overleftrightarrow{D}^\nu \psi_a & a &= \uparrow, \downarrow \\ \psi_{\uparrow, \downarrow} &= \frac{1 \pm \gamma^j \gamma_5}{2} \psi\end{aligned}$$

General parametrization

[Bhoonah, C.L. (2017)]

$$\begin{aligned}\langle p' | \hat{T}_5^{\lambda\mu\nu} | p \rangle &= \bar{u}(p') \left[\frac{P^\nu P^{[\lambda} \Delta^{\mu]} \gamma_5}{2M^2} A_T(t) + \frac{P^\nu P^{[\lambda} \gamma^{\mu]} \gamma_5}{M} B_T(t) \right. \\ &+ \frac{\Delta^\nu \Delta^{[\lambda} \gamma^{\mu]} \gamma_5}{4M} C_T(t) + P^\nu i \sigma^{\lambda\mu} \gamma_5 D_T(t) \\ &\left. + \frac{g^{\nu[\lambda} \Delta^{\mu]} \gamma_5}{2} \tilde{A}_T(t) + M g^{\nu[\lambda} \gamma^{\mu]} \gamma_5 \tilde{B}_T(t) + \frac{P^{[\lambda} i \sigma^{\mu\nu]} \gamma_5}{2} \tilde{D}_T(t) \right] u(p)\end{aligned}$$

Higher twist

Transverse spin-orbit correlations

$$\begin{aligned}C_j &= -\frac{M}{2\sqrt{2}P^+} [B_T(0) + 2\tilde{B}_T(0) + 4\tilde{D}_T(0)] & \text{Asymmetric} \\ \mathbb{C}_j &= -\frac{M}{2\sqrt{2}P^+} [B_T(0) + 2\tilde{B}_T(0) - 2D_T(0)] & \text{Belinfante}\end{aligned}$$

Burkardt's correlation

Instant-form correlation with Belinfante tensor

$$\langle \delta^x J^x \rangle = \mathbb{C}_j|_{\text{IF}}$$

Transversity asymmetry of total AM

Rest frame $\langle \delta^x J^x \rangle = \frac{1}{2} D_T(0)$

[Burkardt (2005)]

[Burkardt (2006)]

Moving frame $\langle \delta^x J^x \rangle = \frac{1}{2} \left[\frac{E-M}{M} B_T(0) + D_T(0) \right]$

[Bhoonah, C.L. (2017)]

Transverse total AM

Rest frame $\langle J^x \rangle = \frac{1}{2} [A(0) + B(0)]$

[Ji (1997)]

[Burkardt (2005)]

Moving frame $\langle J^x \rangle = \frac{1}{2} \left\{ \frac{E-M}{M} B(0) + [A(0) + B(0)] \right\}$

[Leader (2012)]

Link with GPDs

GPD correlators

$$F[\Gamma] = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \Gamma \mathcal{W} \psi(\frac{z^-}{2}) | p \rangle$$

↑
**Straight LF
Wilson line**

Mellin moment of GPDs

$$\frac{1}{2(P^+)^2} \langle p' | \hat{T}^{\mu+} | p \rangle = \int dx x F^{[\gamma^\mu]} \quad \textbf{Vector}$$

$$\frac{1}{2(P^+)^2} \langle p' | \hat{T}_5^{\mu+} | p \rangle = \int dx x F^{[\gamma^\mu \gamma_5]} \quad \textbf{Axial-vector}$$

$$\frac{1}{2(P^+)^2} \langle p' | \hat{T}_5^{\lambda\mu+} | p \rangle = \int dx x F^{[i\sigma^{\lambda\mu} \gamma_5]} \quad \textbf{Tensor}$$

Leading twist ($\mu = +$)

Vector

$$\left\{ \begin{array}{l} \int dx x H(x, \xi, t) = A(t) + 4\xi^2 C(t) \\ \int dx x E(x, \xi, t) = B(t) - 4\xi^2 C(t) \end{array} \right.$$

[Ji (1997)]

Axial-vector

$$\left\{ \begin{array}{l} \int dx x \tilde{H}(x, \xi, t) = \tilde{A}(t) \\ \int dx x \tilde{E}(x, \xi, t) = \tilde{B}(t) \end{array} \right.$$

[Hägler (2004)]
[Diehl, Hägler (2005)]

Tensor

$$\left\{ \begin{array}{l} \int dx x H_T(x, \xi, t) = -\frac{t}{4M^2} A_T(t) - B_T(t) + D_T(t) \\ \int dx x E_T(x, \xi, t) = A_T(t) + B_T(t) \\ \int dx x \tilde{H}_T(x, \xi, t) = -\frac{1}{2} A_T(t) \\ \int dx x \tilde{E}_T(x, \xi, t) = -\xi C_T(t) \end{array} \right.$$

[Hägler (2004)]
[Diehl, Hägler (2005)]

Sub-leading twist

Parametrization [Meissner, Metz, Schlegel (2009)]

Vector

$$\left[\begin{aligned} \int dx x H_{2T}(x, \xi, t) &= 0 \\ \int dx x E_{2T}(x, \xi, t) &= 0 \\ \int dx x \tilde{H}_{2T}(x, \xi, t) &= -2\xi C(t) \\ \int dx x \tilde{E}_{2T}(x, \xi, t) &= -\frac{1}{2} [A(t) + B(t) - D(t)] \end{aligned} \right.$$

[Penttinen *et al.* (2000)]
[Kiptily, Polyakov (2004)]

Axial-vector

$$\left[\begin{aligned} \int dx x H'_{2T}(x, \xi, t) &= \frac{1}{2} [\tilde{A}(t) - \tilde{C}(t) + 2\tilde{F}(t)] + \frac{t}{4M^2} \frac{1}{2} [\tilde{B}(t) - \tilde{D}(t)] \\ \int dx x E'_{2T}(x, \xi, t) &= -\frac{1}{2} [\tilde{A}(t) + \tilde{B}(t) - \tilde{C}(t) - \tilde{D}(t)] \\ \int dx x \tilde{H}'_{2T}(x, \xi, t) &= \frac{1}{4} [\tilde{B}(t) - \tilde{C}(t)] \\ \int dx x \tilde{E}'_{2T}(x, \xi, t) &= 0 \end{aligned} \right.$$

[Penttinen *et al.* (2000)]
[Kiptily, Polyakov (2004)]

Tensor

$$\left[\begin{aligned} \int dx x H'_2(x, \xi, t) &= -\xi \left[\frac{t}{4M^2} C_T(t) + D(t) + \tilde{D}_T(t) \right] \\ \int dx x E'_2(x, \xi, t) &= \xi [C_T(t) + D(t) + \tilde{D}_T(t)] \\ \int dx x \tilde{H}'_2(x, \xi, t) &= -\left[\left(1 - \frac{t}{4M^2}\right) B_T(t) + \tilde{B}(t) - D_T(t) \right] \\ \int dx x \tilde{E}'_2(x, \xi, t) &= \xi \left[\left(1 - \frac{t}{4M^2}\right) A_T(t) + \tilde{A}(t) + D_T(t) \right] \end{aligned} \right.$$

[Bhoonah, C.L. (2017)]

Relations between form factors

Vector

$$\bar{\psi} \gamma^{[\mu} i \overleftrightarrow{D}^{\nu]} \psi = -\varepsilon^{\mu\nu\alpha\beta} \partial_\alpha (\bar{\psi} \gamma_\beta \gamma_5 \psi)$$

[C.L., Mantovani, Pasquini
(in preparation)]

$$D(t) = -G_A(t)$$

Axial-vector

$$\bar{\psi} \gamma^{[\mu} \gamma_5 i \overleftrightarrow{D}^{\nu]} \psi = 2m \bar{\psi} i \sigma^{\mu\nu} \gamma_5 \psi - \varepsilon^{\mu\nu\alpha\beta} \partial_\alpha (\bar{\psi} \gamma_\beta \psi)$$

[C.L. (2014)]

$$\begin{aligned} \tilde{C}(t) &= \frac{m}{2M} H_1(t) - F_1(t) \\ \tilde{D}(t) &= \frac{m}{2M} H_2(t) - F_2(t) \\ \tilde{F}(t) &= \frac{m}{2M} H_3(t) - \frac{1}{2} [F_1(t) + \frac{t}{4M^2} F_2(t)] \end{aligned}$$

Tensor

$$\bar{\psi} i \sigma^{[\lambda\mu} \gamma_5 i \overleftrightarrow{D}^{\nu]} \psi = -2\varepsilon^{\lambda\mu\nu\alpha} \partial_\alpha (\bar{\psi} \psi)$$

[Bhoonah, C.L. (2017)]

$$\bar{\psi} i \sigma^{\lambda\mu} \gamma_5 i \overleftrightarrow{D}_\mu \psi = 2m \bar{\psi} \gamma^\lambda \gamma_5 \psi + i \partial^\lambda (\bar{\psi} \gamma_5 \psi)$$

$$\begin{aligned} D_T(t) + 3\tilde{D}_T(t) &= \Sigma(t) \\ -[(1 - \frac{t}{4M^2}) B_T(t) + 3\tilde{B}_T(t) + \frac{t}{4M^2} C_T(t) - D_T(t)] &= \frac{m}{M} G_A(t) \\ -[(1 - \frac{t}{4M^2}) A_T(t) + 3\tilde{A}_T(t) - C_T(t) + D_T(t)] &= \frac{m}{M} G_P(t) - \Pi(t) \end{aligned}$$

Interpretation of leading-twist relations

Longitudinal OAM

$$L_z = \frac{1}{2} \int dx x [H(x, 0, 0) + E(x, 0, 0)] - \frac{1}{2} G_A(0)$$

$$\rightsquigarrow \langle L_z S_z^N \rangle = \langle J_z S_z^N \rangle - \langle S_z S_z^N \rangle$$

[Ji (1997)]

[Shore, White (2000)]

Longitudinal spin-orbit correlation

$$C_z = \frac{1}{2} \int dx x \tilde{H}(x, 0, 0) - \frac{1}{2} [F_1(0) - \frac{m}{2M} H_1(0)]$$

$$\rightsquigarrow \langle L_z S_z \rangle = \langle J_z S_z \rangle - \langle S_z S_z \rangle$$

[C.L. (2014)]

Transverse spin-orbit correlation

$$\frac{\sqrt{2}P^+}{M} C_j = \frac{1}{3} \int dx x [H_T(x, 0, 0) + \frac{1}{2} \bar{E}_T(x, 0, 0)] - \frac{2}{3} [\Sigma(0) - \frac{m}{2M} G_A(0)]$$

$$\rightsquigarrow \langle L_j T_j \rangle = \langle J_j T_j \rangle - \langle S_j T_j \rangle$$

[Bhoonah, C.L. (2017)]

Some figures

	L_z^u	L_z^d	C_z^u	C_z^d	C_j^u	C_j^d	
$\mu^2 \approx 0.26 \text{ GeV}^2$ {	LFCQM	0.071	0.055	-0.84	-0.54	×	×
	LF χ QSM	-0.008	0.077	-0.80	-0.55	×	×
$\mu^2 = 4 \text{ GeV}^2$	Lattice	-0.175	0.205	-0.90	-0.53	-3.6	-2.2

LFCQM, LF χ QSM

[C.L., Pasquini, Vanderhaeghen (2011)]
[C.L. (2014)]

Lattice QCD

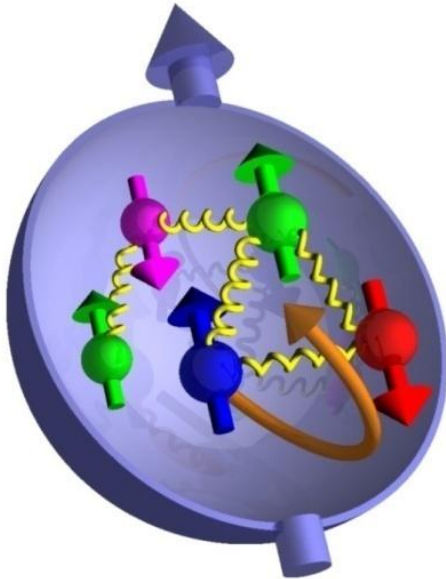
[Göckeler *et al.* (2005)]
[Göckeler *et al.* (2007)]
[Bratt *et al.* (2010)]
[Abdel-Rehim *et al.* (2015)]



Spin and kinetic OAM of valence quarks are **anti-correlated !**

Conclusions

Take home message



- In presence of spin density, EMT is **asymmetric**
- EMT can be decomposed according to parton polarization
- Spin-orbit correlations complementary to nucleon spin sum rule
- Information encoded in GPDs and standard form factors
- Quark models and Lattice suggest **negative** spin-orbit correlations